

# 1) Complete functional Differentiation of the $2p$ orbitals

The problem is to find  $1/r_{vac}$  from:

$$\nabla^2 \psi = - \left( \frac{4\pi c}{\hbar} \right) \frac{1}{r_{vac}} \psi \quad - (1)$$

where:

$$\nabla^2 \psi = \frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial}{\partial \theta} \left( \sin^2 \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \quad - (2)$$

and  $\int_0^{2\pi} d\phi \int_0^\pi \sin \theta d\theta = 4\pi \quad - (3)$

$2p_z$  orbital  
 $\psi(2p_z) = \frac{1}{2} \left( \frac{3}{\pi} \right)^{1/2} \cos \theta \left( \frac{1}{a} \right)^{3/2} \left( \frac{1}{2\sqrt{6}} \right) \frac{r}{a} \exp \left( -\frac{r}{2a} \right) \quad - (4)$

$2p_y$  orbital  
 $\psi(2p_y) = -\frac{1}{2} \left( \frac{3}{\pi} \right)^{1/2} \sin^2 \theta e^{i\phi} \left( \frac{1}{a} \right)^{3/2} \left( \frac{1}{2\sqrt{6}} \right) \frac{r}{a} \exp \left( -\frac{r}{2a} \right) \quad - (5)$

$2p_x$  orbital  
 $\psi(2p_x) = \frac{1}{2} \left( \frac{3}{\pi} \right)^{1/2} \sin^2 \theta e^{-i\phi} \left( \frac{1}{a} \right)^{3/2} \left( \frac{1}{2\sqrt{6}} \right) \frac{r}{a} \exp \left( -\frac{r}{2a} \right) \quad - (6)$

Volume element:  $d\tau = r^2 dr \sin \theta d\theta d\phi \quad - (7)$

Probability of finding the electron in a spherical shell of thickness  $dr$  and radius  $r$  is found by integrating in the range:

$$0 \leq \theta \leq \pi, \quad 0 \leq \phi \leq 2\pi \quad - (8)$$

2) The probability is:

$$P = \int_0^\pi \int_0^{2\pi} |\psi(r, \theta, \phi)|^2 \sin \theta r^2 dr d\phi d\theta$$

— (9)

Method.

Use computer algebra to solve eq. (1) for each orbital, and find  $r(\text{vac})(r, \theta, \phi)$  for each. Then

$$\langle r(\text{vac}) \rangle = \int_0^\pi \int_0^{2\pi} r(\text{vac}) d\phi \sin \theta d\theta, \quad \text{— (10)}$$

for each orbital. The average over three orbitals is:

$$\langle \langle r(\text{vac}) \rangle \rangle_{2p} = \frac{1}{3} \left( \langle r(\text{vac})_x \rangle + \langle r(\text{vac})_y \rangle + \langle r(\text{vac})_z \rangle \right) \quad \text{— (11)}$$

The Lamb shift is proportional to:

$$\frac{1}{\langle r(\text{vac})_{2s} \rangle} - \frac{1}{\langle \langle r(\text{vac}) \rangle \rangle_{2p}} \quad \text{— (12)}$$

$$\text{Thus } \langle r(\text{vac}) \rangle = 4\pi r(\text{vac}) \quad \text{— (13)}$$

This procedure integrates out  $\theta$  and  $\phi$ , i.e.

average over them. Alternatively we could use:

$$\langle \psi^2(2p_z) \rangle = \left( \int_0^\pi \sin \theta \psi_{2p_z}^2 d\theta \right)^{1/2} \quad \text{— (14)}$$

etc.

3) If the average  $\langle r(\text{vac}) \rangle$  vanishes, then it would be possible to write:

$$\langle r^2(\text{vac}) \rangle^{1/2} = \left( \int_0^\pi \int_0^{2\pi} r^2(\text{vac}) d\phi \sin\theta d\theta \right)^{1/2} \quad (15)$$

### Notes

1) In general,  $r(\text{vac})$  is a function of  $r$ ,  $\theta$  and  $\phi$  for each orbital.

2) For the  $2s$  orbital,  $r(\text{vac})$  is a function of only  $r$ . In this case:

$$\langle r(\text{vac})_{2s} \rangle = \int_0^\pi \int_0^{2\pi} r(\text{vac})_{2s} d\phi \sin\theta d\theta = 4\pi r(\text{vac})_{2s}, \quad (16)$$

and:  $\langle r^2(\text{vac})_{2s} \rangle^{1/2} = \langle r(\text{vac})_{2s} \rangle. \quad (17)$

3) The type of averaging is important, e.g.

$$\langle \psi(2p_z) \rangle = \frac{1}{2} \left( \frac{3}{\pi} \right)^{1/2} \left( \frac{1}{a} \right)^{3/2} \left( \frac{1}{2\sqrt{6}} \right) \frac{r}{a} \exp\left(-\frac{r}{2a}\right) \int_0^\pi \sin\theta \cos\theta d\theta$$

— (18)

So this type of averaging cannot be used.