

60(4). Magnetostatics is Generally Covariant Unified Field Theory, Resonance Effect, Gravitomagnetics

The ECE magnetostatic equations are:

$$\underline{B}^a = \underline{\nabla} \times \underline{A}^a - \underline{\omega}^a{}_b \times \underline{A}^b \quad - (1)$$

$$\underline{\nabla} \cdot \underline{B}^a = \mu_0 \tilde{j}^{a0} \quad - (2)$$

$$\underline{\nabla} \times \underline{B}^a = \frac{\mu_0}{c} \underline{j}^a \quad - (3)$$

i.e. $\underline{\nabla} \cdot (\underline{\omega}^a{}_b \times \underline{A}^b) = -\mu_0 \tilde{j}^{a0} \quad - (4)$

$$\underline{\nabla} \times (\underline{\nabla} \times \underline{A}^a - \underline{\omega}^a{}_b \times \underline{A}^b) = \frac{\mu_0}{c} \underline{j}^a \quad - (5)$$

Eq. (4) is the generally covariant Gauss Law and

Eq. (5) is the generally covariant Ampère Law. In

these equations:

$$\underline{\omega}^a{}_b \times \underline{A}^b = \underline{\omega}^a{}_1 \times \underline{A}^1 + \underline{\omega}^a{}_2 \times \underline{A}^2 + \underline{\omega}^a{}_3 \times \underline{A}^3 \quad - (6)$$

If: $\tilde{j}^{a0} = \tilde{j}_1^{a0} + \tilde{j}_2^{a0} + \tilde{j}_3^{a0} \quad - (7)$

equation (4) for example splits into three equations:

$$\underline{\nabla} \cdot (\underline{\omega}^a{}_1 \times \underline{A}^1) = -\mu_0 \tilde{j}_1^{a0} \quad - (8)$$

$$\underline{\nabla} \cdot (\underline{\omega}^a{}_2 \times \underline{A}^2) = -\mu_0 \tilde{j}_2^{a0} \quad - (9)$$

$$\underline{\nabla} \cdot (\underline{\omega}^a{}_3 \times \underline{A}^3) = -\mu_0 \tilde{j}_3^{a0} \quad - (10)$$

2) So the rotation can be simplified by considering each equation to be of the form:

$$\underline{\nabla} \cdot (\underline{\omega} \times \underline{A}) = -\mu_0 \tilde{j}^0 \quad (11)$$

i.e.

$$\underline{\omega} \times \underline{A} = -\mu_0 \int \tilde{j}^0 d\underline{r} \quad (12)$$

Similarly, eq. (5) can be split into ~~two~~ equations of the form:

$$\underline{\nabla} \times (\underline{\nabla} \times \underline{A} - \underline{\omega} \times \underline{A}) = \frac{\mu_0}{c} \underline{\tilde{J}} \quad (13)$$

so:

$$\underline{\nabla} \times (\underline{\nabla} \times \underline{A}) = \frac{\mu_0}{c} \underline{\tilde{J}} - \mu_0 \underline{\nabla} \times \int \tilde{j}^0 d\underline{r} \quad (14)$$

Using vector identities, eq. (13) is:

$$\begin{aligned} \underline{\nabla} (\underline{\nabla} \cdot \underline{A}) - \nabla^2 \underline{A} - \underline{\omega} (\underline{\nabla} \cdot \underline{A}) + (\underline{\omega} \cdot \underline{\nabla}) \underline{A} + (\underline{\nabla} \cdot \underline{\omega}) \underline{A} - (\underline{A} \cdot \underline{\nabla}) \underline{\omega} \\ = \frac{\mu_0}{c} \underline{\tilde{J}} \end{aligned} \quad (15)$$

Therefore \underline{A} is amplified at resonance by the current density $\underline{\tilde{J}}$.

3) Consider the Z component of $\underline{\tilde{J}}$:

$$\frac{\partial^2 A_x}{\partial z \partial x} + \frac{\partial^2 A_y}{\partial z \partial y} + \omega_x \frac{\partial A_z}{\partial x} + \omega_y \frac{\partial A_z}{\partial y} - \omega_z \frac{\partial A_x}{\partial x} - \omega_z \frac{\partial A_y}{\partial y} + \left(\frac{\partial \omega_x}{\partial x} + \frac{\partial \omega_y}{\partial y} \right) A_z - \frac{\partial \omega_z}{\partial x} A_x - \frac{\partial \omega_z}{\partial y} A_y = \frac{\mu_0}{c} \tilde{J}_z$$

- (16)

Finally assume that only A_x is non-zero:

$$\frac{\partial}{\partial z} \left(\frac{\partial A_x}{\partial x} \right) - \omega_z \frac{\partial A_x}{\partial x} - \left(\frac{\partial \omega_z}{\partial x} \right) A_x = \frac{\mu_0}{c} \tilde{J}_z$$

- (17)

This is a relatively simple resonance equation producing a surge of A_x at resonance. Here A_x must be a function of both x and z and ω_z must be a function of x .

So at resonance, amplification of the magnetic field occurs by virtue of the spin connection. Gravitomagnetic effects are described by eq. (14) via \tilde{j} , and by eq. (17) via \tilde{J}_z .

Gravitomagnetic effects are amplified at resonance.