

Notes for Paper 55, Part One

Derivation of the Euler Equations from ECE

Define the angular momentum as:

$$\underline{J}^a = \underline{J}^{(0)} \underline{q}^a \quad - (1)$$

(pp 287 ff of volume one). Now we use the canonical structure equations:

$$\underline{T}^a = d \wedge \underline{q}^a + \omega^a_b \wedge \underline{q}^b \quad - (2)$$

to find the torque:

$$\underline{N}^a = c (d \wedge \underline{J}^a + \omega^a_b \wedge \underline{J}^b) \quad - (3)$$

This is a generally covariant Euler equation of motion, more accurately the g.c. rotational equation.

In the classical limit the Euler equation of motion is:

$$\begin{aligned} \underline{N} &= \left(\frac{d\underline{J}}{dt} \right)_{\text{lab frame}} \\ &= \left(\frac{d\underline{J}}{dt} \right)_{\text{moving frame}} + \underline{\omega} \times \underline{J} \end{aligned}$$

2) As described by Marion and Thornton on page 388 ff:

$$\underline{N} = \left(\frac{dL}{dt} \right)_{\text{fixed}} \quad - (5)$$

is valid in an inertial frame only.

Eq. (3) is automatically relativistic, and shows that:

$$\boxed{N^a = c J^{(0)} T^a} \quad - (6)$$

relating the torque from N^a and the Carter torsion T^a . The Carter structure equation defines the torque in general relativity. The angular momentum is a type of tetrad field, a fact inferred in vol. one. In eq. (4) $\underline{\omega}$ is the angular velocity in the lab frame in classical dynamics.

Eq. (3) is closely analogous to the equation defining the electromagnetic field:

$$F^a = d \wedge A^a + \omega^a{}_b \wedge A^b$$

3) So the $\omega^a{}_b \wedge J^b$ term originates in a spinning frame of reference, i.e. spinning spacetime itself. If spacetime is not spinning there is no Cartan torsion, so there is no torque:

$$N^a = 0. \quad - (6)$$

This situation describes Newtonian dynamics, which is the weak field limit of a Riemann geometry which has curvature $R^a{}_b$, but no torsion, T^a . In Newtonian dynamics the torque is well known to be:

$$\underline{N} = \underline{r} \times \underline{F} = \underline{\dot{J}}, \quad - (7)$$

where \underline{F} is force and \underline{r} is arm length. The Newtonian force \underline{F} originates in the Riemann curvature and not in the Cartan torsion.

So in classical dynamics there is a mixture of concepts and an internal inconsistency, classical dynamics attempts to describe torque without torsion. It is ^{self} consistent and generally covariant description of rotational

4) dynamics is eq. (3). Therefore there are rotational effects in EFE but do not exist in classical dynamics. From eq. (6) it is found that torque is the Carter torsion with a factor $c J^{(0)}$. So torque interacts with a gravitational field through the first Bianchi identity:

$$\begin{aligned}
 D \wedge N^a + \omega^a_b \wedge N^b & \\
 &= c J^{(0)} R^a_b \wedge v^b \\
 &= c R^a_b \wedge J^b,
 \end{aligned}
 \tag{8}$$

i.e.

$$D \wedge N^a = c R^a_b \wedge J^b \tag{9}$$

In the absence of a gravitational field:

$$R^a_b = 0 \tag{10}$$

$$\text{so: } D \wedge N^a = 0 \tag{11}$$

The Newtonian torque (7) is calculated from:

$$D \wedge R^a_b = 0 \tag{12}$$