

or(2): Relativistic Quantum Mechanics and the Thomas Half  
The hamiltonian of the Sommerfeld and Dirac atom can be  
given as:

$$H = \gamma mc^2 + U - (1)$$

$\gamma$  is defined by the Thomas half:

$$\frac{\Delta\phi_T}{2\pi} = \gamma - 1 - (2)$$

$$\text{so } H = \left(1 + \frac{\Delta\phi_T}{2\pi}\right)mc^2 + U - (3)$$

$$\text{It follows that: } H_0 = H - mc^2 = \frac{\Delta\phi_T}{2\pi} mc^2 + U - (4)$$

$$= (\gamma - 1)mc^2 + U$$

$$:= T + U$$

in which the relativistic kinetic energy is:

$$T = (\gamma - 1)mc^2 = \frac{\Delta\phi_T}{2\pi} mc^2 - (5)$$

The non-relativistic kinetic energy is defined by:

$$T_0 = \left( \left(1 - \frac{v_N^2}{c^2}\right)^{-1/2} - 1 \right) mc^2$$

$$\xrightarrow{v_N \ll c} \left(1 + \frac{1}{2} \frac{v_N^2}{c^2} - 1\right) mc^2 - (6)$$

$$= \frac{1}{2} m v_N^2$$

$$\text{so } T_0 = \frac{\Delta\phi_T}{2\pi} mc^2 = \frac{1}{2} \frac{v_N^2}{c^2} mc^2 - (7)$$

Both relativistic and non-relativistic kinetic energy are defined by the Thomas half. So the energy levels of all atoms and molecules are defined by:

$$\langle H_0 \rangle = mc^2 \left\langle \frac{\Delta \phi_T}{2\pi} \right\rangle + \langle u \rangle - (8)$$

So all relativistic and non-relativistic levels. The transition from non-relativistic to relativistic atom is defined by:

$$\frac{1}{2} \frac{\sqrt{N}}{c^2} \rightarrow \gamma - 1 - (9)$$

In Q. Schrödinger atom:

$$\left\langle \frac{\sqrt{N}}{c^2} \right\rangle = \frac{d}{n^2} - (10)$$

and

$$\langle u \rangle = - \frac{d}{n^2} mc^2 - (11)$$

where  $d$  is the fine structure constant and  $n$  is the principal quantum number.

So from eq. (8):

$$\begin{aligned} \langle H_0 \rangle &= mc^2 \left\langle \frac{\Delta \phi_T}{2\pi} \right\rangle - \frac{d}{n^2} mc^2 \\ &= \frac{1}{2} \left\langle \frac{\sqrt{N}}{c^2} \right\rangle mc^2 - \frac{d}{n^2} mc^2 \\ &= \frac{1}{2} \frac{d}{n^2} mc^2 - \frac{d}{n^2} mc^2 \\ &= - \frac{1}{2} \frac{d}{n^2} mc^2 - (12) \end{aligned}$$

The expectation value of the Thomas half is the

Schrodinger H atom is:

$$\left\langle \frac{\Delta \phi}{2\pi} \right\rangle = \frac{1}{2} \frac{\vec{d}^2}{\vec{n}^2} - (13)$$

$$= \frac{1}{2} \left\langle \frac{\vec{v}_n^2}{c^2} \right\rangle$$

so

$$\left\langle \frac{\vec{v}_n^2}{c^2} \right\rangle = \frac{\vec{d}^2}{\vec{n}^2} - (14)$$

It follows that

$$\left\langle \frac{\vec{p}^2}{2m} \right\rangle = mc^2 \frac{\vec{d}^2}{\vec{n}^2} - (15)$$

The expectation value is defined by:

$$\left\langle \frac{\vec{f}^2}{2m} \right\rangle = \int \psi^* \frac{\vec{f}^2}{2m} \psi d\tau - (16)$$

$$= -\frac{\vec{f}^2}{2m} \int \psi^* \vec{J}^2 \psi d\tau$$

The wavefunctions are calculated from:

$$\langle H_0 \rangle := E = -\frac{\vec{f}^2}{2m} \vec{J}^2 \psi - \frac{e^2}{4\pi\epsilon_0 r} \psi - (17)$$

where

$$\begin{aligned} \vec{u} &= -\frac{e}{4\pi\epsilon_0 r} - (18) \\ &= -\frac{d \vec{h} c}{r} \end{aligned}$$

Final expression:

$$E^2 = \gamma^2 \vec{n}^2 c^4 = \vec{c}^2 \vec{p}^2 + \vec{n}^2 c^4 - (19)$$

it follows that:

$$4) H = \gamma mc^2 + U \\ = \left( c_p^2 + m^2 c^4 \right)^{1/2} + U - (20)$$

so  $(H-U) = c_p^2 + m^2 c^4 - (21)$

The relativistic kinetic energy is :

$$T = (\gamma - 1) mc^2 = \frac{\Delta \phi_T}{2\pi} mc^2 - (22)$$

The total relativistic energy is :

$$E = \gamma mc^2 = T + mc^2 - (23)$$

so  $E = H - U - (24)$

From eq. (21) :

$$(E - mc^2)(E + mc^2) = p^2 c^2 - (24)$$

and

$$E = \frac{p^2 c^2}{H - U + mc^2} + mc^2 - (25)$$

Therefore

$$H = E + U = \frac{p^2 c^2}{H - U + mc^2} + mc^2 + U - (26)$$

and

$$H_0 = H - mc^2 = \frac{p^2 c^2}{H - U + mc^2} + U - (27)$$

Comparing eq.s. (4) and (27) :

$$\boxed{\frac{\Delta \phi_T}{2\pi} = (\gamma - 1) = \frac{1}{mc^2} \frac{p^2 c^2}{H - U + mc^2}} - (28)$$

Eq. (28) is the route to the quantization of:

$$H_0 = (V-1)mc^2 + U - (29)$$
$$= \frac{p^2 c^2}{2m} + U$$

So the Thomas half gives the energy levels of the Dirac equation, spin-orbit coupling, and the electron's factor.

The usual route to quantization is to use the Dirac approximation of the previous note:

$$H_0 = \frac{p^2}{2m} \left( 1 + \frac{U}{2mc^2} \right) + U - (30)$$

so

$$H_0 \psi = -\frac{\hbar^2}{2m} \nabla^2 \psi + \frac{1}{4mc^2} \sigma \cdot p U \psi - (31)$$

$$= -\frac{\hbar^2}{2m} \nabla^2 \psi + U \psi - i \frac{\hbar}{4mc^2} \sigma \cdot \nabla (U \sigma \cdot p \psi)$$

giving the spin-orbit correction to the energy levels of the atom.

The spin-orbit fine structure of  $H$  is due entirely to the Thomas half defined in Eq. (28). In the Schrödinger case:

$$H_0 \psi = -\frac{\hbar^2}{2m} \nabla^2 \psi + U \psi - (32)$$

$$\langle H_0 \rangle = -\frac{1}{2} \frac{d^2}{n^2} mc^2 - (33)$$

The fine structure is superimposed on Eq. (33) and is calculated accurately by Eq. (31). In addition, the Lamb shift appears in which  $H$  and is due to the effect of the

<sup>6)</sup> vacuum. In ECE2 unified field theory the Lamb shift is described by the following equation:

$$F = \frac{d}{dt} (\gamma m \dot{\phi}) = -\frac{dp c}{mc} + \omega \phi - (34)$$

where  $\omega$  is the spin correction. The latter is found for the experimentally observed Lamb shift.

In some previous UFT papers eq. (29) has been expressed as:

$$H_0 = \frac{p^2 c^2}{T + mc^2} + U - (35)$$

$$= \frac{p^2 c^2}{(\gamma + 1)mc^2} + U$$

$$= \frac{p^2}{m(\gamma + 1)} + U$$

which can now be written as:

$$H_0 = \frac{p^2}{m \left( \frac{\Delta \phi_T}{2\pi} + 2 \right)} + U - (36)$$

The classical result:

$$H_0 = \frac{p^2}{2m} + U - (37)$$

is obtained when

$$\frac{\Delta \phi_T}{2\pi} \rightarrow 0 - (38)$$