

394(3): Conservation of Antisymmetry in the Macroscopic Zitterbewegung Theory

The Riemannian torsion tensor $T^\lambda_{\mu\nu}$ is defined by the well known commutator equation:

$$[D_\mu, D_\nu]V^\rho = -T^\lambda_{\mu\nu} D_\lambda V^\rho + R^\rho_{\sigma\mu\nu} V^\sigma \quad (1)$$

where $T^\lambda_{\mu\nu} = \Gamma^\lambda_{\mu\nu} - \Gamma^\lambda_{\nu\mu} \quad (2)$

and where $\Gamma^\lambda_{\mu\nu}$ is the Christoffel connection. This must be antisymmetric:

$$\Gamma^\lambda_{\mu\nu} = -\Gamma^\lambda_{\nu\mu} \quad (3)$$

because if it were symmetric:

$$\mu = ? \nu \quad (4)$$

and $[D_\mu, D_\nu] = ? 0 \quad (5)$

Eq. (3) alone is enough to refute Einstein's general relativity.

Using the tetrad postulate:

$$\Gamma^a_{\mu\nu} = \delta^a_\mu \delta^b_\nu + \omega^a_{\mu b} \delta^b_\nu \quad (6)$$

it follows that:

$$\delta^a_\mu \delta^b_\nu + \omega^a_{\mu b} \delta^b_\nu = -(\delta^a_\nu \delta^b_\mu + \omega^a_{\nu b} \delta^b_\mu) \quad (7)$$

Eq. (7) gives the scalar and vector antisymmetry laws. They are a direct consequence of eq. (1), which defines the Minkowski Curvature structure equations and the Riemann curvature and torsion tensors.

2) The Lindstrom or trace antisymmetry law is also a direct consequence of fundamental geometry:

$$\Gamma_{00}^a + \Gamma_{11}^a + \Gamma_{22}^a + \Gamma_{33}^a = 0. \quad (8)$$

It is also true that:

$$\Gamma_{00}^a = \Gamma_{11}^a = \Gamma_{22}^a = \Gamma_{33}^a = 0. \quad (9)$$

Now define:

$$\begin{aligned} \omega_{\mu\nu}^a q_\nu^b &= \omega_{\mu 0}^a q_\nu^0 + \omega_{\mu 1}^a q_\nu^1 + \omega_{\mu 2}^a q_\nu^2 + \omega_{\mu 3}^a q_\nu^3 \\ &= \omega_{\mu\nu}^a q_\nu^b \end{aligned} \quad (10)$$

The Cartan torsion tensor:

$$T_{\mu\nu}^a = \partial_\mu q_\nu^a - \partial_\nu q_\mu^a + \omega_{\mu b}^a q_\nu^b - \omega_{\nu b}^a q_\mu^b \quad (11)$$

reduces for each a to:

$$T_{\mu\nu}^a = (\partial_\mu + \omega_\mu^a) q_\nu^a - (\partial_\nu + \omega_\nu^a) q_\mu^a. \quad (12)$$

The antisymmetry laws become:

$$(\partial_\mu + \omega_\mu^a) q_\nu^a = -(\partial_\nu + \omega_\nu^a) q_\mu^a \quad (13)$$

and:

$$\Gamma_{00}^a = \partial_0 q_0^a + \omega_0^a q_0^a = 0 \quad (14)$$

$$\Gamma_{ii}^a = (\partial_i + \omega_i^a) q_i^a = 0 \quad (15)$$

$$i = 1, 2, 3.$$

The antisymmetry laws are a direct consequence of fundamental geometry. They are true for all of physics.

To convert to vector notation:

$$d_\mu = \left(\frac{1}{c} \frac{d}{dt}, \underline{\nabla} \right) \quad - (16)$$

$$q_\mu = \left(q_0, -\underline{q} \right) \quad - (17)$$

Trace Anisotropy Law

$$\frac{dq_0}{dt} + \omega_0 q_0 = 0 \quad - (18)$$

$$\underline{\nabla} \cdot \underline{q} + \underline{\omega} \cdot \underline{q} = 0 \quad - (19)$$

$$\omega_\mu = \left(\frac{\omega_0}{c}, -\underline{\omega} \right) \quad - (20)$$

using

$$\frac{dq_x}{dx} + \omega_x q_x = 0 \quad - (21)$$

Also:

$$\frac{dq_y}{dy} + \omega_y q_y = 0 \quad - (22)$$

$$\frac{dq_z}{dz} + \omega_z q_z = 0 \quad - (23)$$

Scalar Anisotropy Law

$$(d_0 + \omega_0) q_0 = - (d_0 + \omega_0) q_0 \quad - (24)$$

Vector Anisotropy Law

$$(d_\mu + \omega_\mu) q_\nu = - (d_\nu + \omega_\nu) q_\mu \quad - (25)$$

$$\mu, \nu = 1, 2, 3$$

In vector notation eq. (24) is:

$$-\underline{\nabla} \varphi_0 + \underline{\omega} \varphi_0 = -\frac{1}{c} \frac{\partial \varphi}{\partial t} + \underline{\omega}_0 \varphi - (26)$$

Application to Electrodynamics

The ECE potential is used:

$$A_\mu^a = A^{(a)} \varphi_\mu^a - (27)$$

$$\underline{E} = -\underline{\nabla} \phi + \underline{\omega} \phi = -\frac{\partial A}{\partial t} - \underline{\omega}_0 \underline{A} - (28)$$

$$\underline{B} = \underline{\nabla} \times \underline{A} - \underline{\omega} \times \underline{A} - (29)$$

The vector antisymmetry equations are:

$$\frac{\partial A_z}{\partial t} + \frac{\partial A_y}{\partial z} = \omega_y A_z + \omega_z A_y - (30)$$

$$\frac{\partial A_x}{\partial z} + \frac{\partial A_z}{\partial x} = \omega_z A_x + \omega_x A_z - (31)$$

$$\frac{\partial A_y}{\partial x} + \frac{\partial A_x}{\partial y} = \omega_x A_y + \omega_y A_x - (32)$$

Trace antisymmetry means that:

$$\frac{\partial \phi}{\partial t} + \omega_0 \phi = 0 - (33)$$

$$\frac{\partial A_x}{\partial x} + \omega_x A_x = 0 - (34)$$

$$\frac{\partial A_y}{\partial y} + \omega_y A_y = 0 - (35)$$

$$\frac{\partial A_z}{\partial z} + \omega_z A_z = 0 - (36)$$

where

$$A_\mu = \left(\frac{\phi}{c}, -\underline{A} \right) - (37)$$

Electrostatics

$$\underline{E} = -\underline{\nabla} \phi + \underline{\omega} \phi = -\frac{\partial \underline{A}_E}{\partial t} - \omega_0 \underline{A}_E \quad (38)$$

$$\frac{\partial \phi}{\partial t} + \omega_0 \phi = 0 \quad (39)$$

and $\underline{A}_B = \underline{0} \quad (40)$

In the MZ theory, the macroscopic electromagnetic theory:

$$\underline{E} = \underline{E}_0 + \underline{\omega} \phi_0 \quad (41)$$

where \underline{E}_0 and ϕ_0 are the electric field strength and potential in the hypothetical absence of vacuum. If the Coulomb potential or dipole potential is used,

$$\frac{\partial \phi}{\partial t} = 0 \quad (42)$$

$$\omega_0 = 0 \quad (43)$$

so

It follows that:

$$\underline{E} = -\frac{\partial \underline{A}_E}{\partial t} \quad (44)$$

Eqns (34) to (38) reduce to zero = zero.

Magnetostatics

In the MZ theory the magnetic flux density in the presence of vacuum is:

$$\underline{B} = \underline{B}_0 - \underline{\omega} \times \underline{A}_B \quad (45)$$

where

$$\underline{B}_0 = \underline{\nabla} \times \underline{A}_B \quad (46)$$

In magnetostatics:

b)

$$\phi = 0, \quad \underline{A} \cdot \underline{E} = 0 \quad - (47)$$

In the MZ theory, \underline{B} , \underline{B}_0 and $\underline{\omega}$ are known for examples such as the dipole field. So the $\underline{\omega}_i$ symmetry laws (30) to (32) are true automatically in 2T magnetostatics, because of the special connection vector $\underline{\omega}$ is derived from \underline{B} , \underline{B}_0 and \underline{A}_B .

Adding eqns. (34) to (36) gives the gauge equation:

$$\nabla \cdot \underline{A}_B + \underline{\omega} \cdot \underline{A}_B = 0 \quad - (48)$$

In the MZ theory, $\underline{\omega}$ and \underline{A}_B are known, so $\nabla \cdot \underline{A}_B$ can be found, together with $\partial A_{Bx} / \partial x$, $\partial A_{By} / \partial y$, and $\partial A_{Bz} / \partial z$.

The self consistent procedure is to find ω_x , ω_y and ω_z from eqns. (34) to (36) for a given magnetic potential such as the dipole potential. The shivering magnetic field is the inverse of the vacuum is then found from:

$$\underline{B} = \nabla \times \underline{A} - \underline{\omega} \times \underline{A} \quad - (49)$$

and the non square fluctuations in \underline{B} determined, giving $\langle \underline{\delta r} \cdot \underline{\delta r} \rangle$. In tensor notation eqn. (49)

$$B_{ij} = \partial_i A_j - \partial_j A_i + \omega_i A_j - \omega_j A_i \quad - (50)$$

The vector antisymmetry law is:

$$B_{ij} = -B_{ji} \quad - (51)$$

7) So:
$$j_i A_j + \omega_i A_i = - (j_j A_i + \omega_j A_i) - (52)$$
 is true automatically. So eqs. (30) to (32) are true automatically, QED.

It is well emphasizing that $U(1)$ gauge theory is rejected completely in UFT131, using the standard model's own definitions:

$$\begin{aligned} [D_\mu, D_\nu] \phi &= [d_\mu - ig A_\mu, d_\nu - ig A_\nu] \phi \\ &= -ig ([d_\mu, A_\nu] - ig [A_\mu, A_\nu]) \phi \\ &= -ig [d_\mu, A_\nu] \phi \\ &= -ig (d_\mu (A_\nu \phi) - A_\nu d_\mu \phi) \\ &= -ig ((d_\mu A_\nu) \phi + A_\nu (d_\mu \phi) - A_\nu (d_\mu \phi)) \\ &= -ig d_\mu A_\nu \phi \quad - (53) \end{aligned}$$

It follows that in $U(1)$ gauge theory:

$$d_\mu A_\nu = - d_\nu A_\mu \quad - (54)$$

but as shown in UFT131 and UFT132 this is not true, so the entire standard unified field theory is, really, rejected.
