

$\beta(\vec{r})$: Ensemble Average for Dipole Potential

To first order in x :

$$\langle \phi \rangle = \frac{1}{4\pi\epsilon_0 r^3} \left\langle (\underline{r} + \delta\underline{r}) \cdot \underline{p} \left(1 - \frac{3x}{2} \right) \right\rangle - (1)$$

$$\text{let } x = \frac{1}{r^2} (2\underline{r} \cdot \delta\underline{r} + \delta\underline{r} \cdot \delta\underline{r}). - (2)$$

then

$$\text{So: } \langle \phi \rangle = \frac{1}{4\pi\epsilon_0 r^3} \left(\underline{r} \cdot \underline{p} + \langle \delta\underline{r} \cdot \underline{p} \rangle \right)$$

$$= \frac{3\underline{p}}{8\pi\epsilon_0 r^5} \cdot \left\langle (\underline{r} + \delta\underline{r})(2\underline{r} \cdot \delta\underline{r} + \delta\underline{r} \cdot \delta\underline{r}) \right\rangle - (3)$$

$$\text{By isotropy: } \langle \delta\underline{r} \cdot \underline{p} \rangle = 0 - (4)$$

In eq. (3):

$$\left\langle (\underline{r} + \delta\underline{r})(2\underline{r} \cdot \delta\underline{r} + \delta\underline{r} \cdot \delta\underline{r}) \right\rangle$$

$$\left\langle (\underline{r} + \delta\underline{r})(2\underline{x} \cdot \delta\underline{x} + 2\underline{y} \cdot \delta\underline{y} + 2\underline{z} \cdot \delta\underline{z}) \right\rangle - (5)$$

$$= \frac{1}{2} \left\langle ((\underline{x} + \delta\underline{x})(2\underline{x} \cdot \delta\underline{x} + 2\underline{y} \cdot \delta\underline{y} + 2\underline{z} \cdot \delta\underline{z}) + \dots + \underline{\delta x}^2 + \underline{\delta y}^2 + \underline{\delta z}^2) \right\rangle + \dots$$

$$\text{By isotropy: } \langle \underline{x} \cdot \delta\underline{x} \rangle = \langle \underline{y} \cdot \delta\underline{y} \rangle = \langle \underline{z} \cdot \delta\underline{z} \rangle = 0 - (6)$$

$$\text{so } \left\langle (\underline{r} + \delta\underline{r})(2\underline{r} \cdot \delta\underline{r} + \delta\underline{r} \cdot \delta\underline{r}) \right\rangle - (7)$$

$$= 3\underline{r} \cdot \langle \delta\underline{r} \cdot \delta\underline{r} \rangle$$

Therefore to first order in x :

$$\langle \phi \rangle = \frac{\underline{P} \cdot \underline{E}}{4\pi \epsilon_0 r^3} \left(1 - \frac{q}{2r} \langle \underline{s}_r \cdot \underline{s}_r \rangle \right) - (8)$$

The standard molecular dipole potential is:

$$\phi = \frac{\underline{P} \cdot \underline{E}}{4\pi \epsilon_0 r^3} - (9)$$

As in the case of the Landau shift in H, the vacuum contributes a term $\langle \underline{s}_r \cdot \underline{s}_r \rangle$. The true dipole potential is always eq. (8) to first order in \propto . This is because the vacuum is always present. This means that the radiative corrections are always small, and exclusively dominate the fact that the vacuum creates potential and field in matter and in circuits.

The vacuum charge \propto dependence of the dipole potential. To first order in \propto it becomes the sum of terms in $1/r^3$ and $1/r^5$. The ensemble averaged electric field strength due to $\langle \phi \rangle$ will be calculated in the next note to first order in \propto .