

1(2). The Spin Cyclic Theorem
 The well known $B^{(3)}$ field of ECE2 unified field theory is defined by:

$$\underline{B}^{(3)*} = -\frac{i\kappa}{A^{(0)}} \underline{A} \times \underline{A}^* \\ = -\frac{i\kappa}{A^{(0)}} \underline{A}^{(1)} \times \underline{A}^{(2)} \quad - (1)$$

where

$$\underline{A} \times \underline{A}^* = \underline{A}^{(1)} \times \underline{A}^{(2)} \quad - (2)$$

is the cross product of vector potentials, notably plane waves. For example, if:

$$\underline{A}^{(1)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i(\omega t - \kappa z)} \quad - (3)$$

$$\underline{A}^{(2)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i(\omega t - \kappa z)} \quad - (4)$$

$$\underline{A}^{(1)} \times \underline{A}^{(2)} = \frac{A^{(0)2}}{2} \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ 1 & -i & 0 \\ 1 & i & 0 \end{vmatrix} \quad - (5) \\ = i A^{(0)2} \underline{k}$$

$$\underline{B}^{(3)*} = \kappa A^{(0)} \underline{k} \\ = B^{(0)} \underline{k} \quad - (6) \\ = B^{(0)} \underline{e}^{(3)*}$$

The B Cyclic Theorem follows:

$$\underline{B}^{(1)} \times \underline{B}^{(2)} = i \underline{B}^{(0)} \underline{B}^{(3)*} - (7)$$

$$\underline{B}^{(3)} \times \underline{B}^{(1)} = i \underline{B}^{(0)} \underline{B}^{(2)*} - (8)$$

$$\underline{B}^{(2)} \times \underline{B}^{(3)} = i \underline{B}^{(0)} \underline{B}^{(1)*} - (9)$$

in the complex circular basis:

$$\underline{e}^{(1)} \times \underline{e}^{(2)} = i \underline{e}^{(3)*} - (10)$$

$$\underline{e}^{(3)} \times \underline{e}^{(1)} = i \underline{e}^{(2)*} - (11)$$

$$\underline{e}^{(2)} \times \underline{e}^{(3)} = i \underline{e}^{(1)*} - (12)$$

Here

$$\underline{e}^{(1)} = \frac{1}{\sqrt{2}} (\underline{i} - i \underline{j}) - (13)$$

$$\underline{e}^{(2)} = \frac{1}{\sqrt{2}} (\underline{i} + i \underline{j}) - (14)$$

$$\underline{e}^{(3)} = \underline{k} - (15)$$

where the Cartesian basis is:

$$\underline{i} \times \underline{j} = \underline{k} - (16)$$

$$\underline{k} \times \underline{i} = \underline{j} - (17)$$

$$\underline{j} \times \underline{k} = \underline{i} - (18)$$

In ECE2 electrodynamics:

$$\underline{B} = \nabla \times \underline{A} - \underline{\omega} \times \underline{A} - (19)$$

where $\underline{\omega}$ is the spin angular velocity vector. Here.

$$\underline{B} (\text{interaction w/ vacuum}) = - \underline{\omega} \times \underline{A} - (20)$$

It follows that:

$$\underline{B}^{(3)*} = -i \underline{B}^{(1)} \times \underline{B}^{(2)} = -i \frac{\kappa}{A^{(0)}} \underline{A}^{(1)} \times \underline{A}^{(2)}$$

$$= -i \underline{\omega}^{(1)} \times \underline{A}^{(2)} \quad - (21)$$

s. the spin current is:

$$\underline{\omega}^{(1)} = \frac{\kappa}{A^{(0)}} \underline{A}^{(1)} \quad - (22)$$

Similarly:

$$\underline{\omega}^{(2)} = \frac{\kappa}{A^{(0)}} \underline{A}^{(2)} \quad - (23)$$

It follows that:

$$\underline{B}^{(3)*} = -i \frac{A^{(0)}}{\kappa} \underline{\omega}^{(1)} \times \underline{\omega}^{(2)} \quad - (24)$$

$$= -i \frac{B^{(0)}}{\kappa} \underline{\omega}^{(1)} \times \underline{\omega}^{(2)}$$

Note carefully that

$$\underline{A}^{(3)} = \underline{\omega}^{(3)} = \underline{0} \quad - (25)$$

because polar vectors cannot be obtained from a conjugate product by symmetry.

Now define the magnetization of the vacuum

by:

$$4) \underline{B}^{(3)*} = \frac{1}{\mu_0} \underline{M}^{(3)*} = -\frac{i\hbar}{A^{(0)}} \underline{A} \times \underline{A}^* = -i\omega \times \underline{A}^* - (26)$$

where
$$\underline{\omega} = \frac{\omega^{(0)}}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i\phi} - (27)$$

and
$$\underline{A}^* = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} - (28)$$

In eq. (26):

$$\underline{\omega} \times \underline{A}^* = \frac{\omega^{(0)} A^{(0)}}{2} \begin{vmatrix} \underline{i} & \underline{j} & \underline{k} \\ 1 & -i & 0 \\ 1 & i & 0 \end{vmatrix} - (29)$$

$$= i\omega^{(0)} A^{(0)} \underline{k}$$

Note that:
$$\omega_x A_y^* = -\omega_y A_x^* - (30)$$

and
$$\omega_z = A_z^* = 0 - (31)$$

Eqs. (30) and (31) are solutions of the

antisymmetry equations:

$$\frac{\partial A_z^*}{\partial y} + \frac{\partial A_y^*}{\partial z} = \omega_y A_z^* + \omega_z A_y^* - (32)$$

$$\frac{\partial A_x^*}{\partial z} + \frac{\partial A_z^*}{\partial x} = \omega_z A_x^* + \omega_x A_z^* - (33)$$

$$\frac{\partial A_y^*}{\partial x} + \frac{\partial A_x^*}{\partial y} = \omega_x A_y^* + \omega_y A_x^* - (34)$$

Eq. (32) to (34) reduce to:

$$\frac{\partial A_y^*}{\partial z} = \omega_z A^* - (35)$$

$$\frac{\partial A_x^*}{\partial z} = \omega_z A_x^* - (36)$$

$$\omega_x A_y^* = -\omega_y A_x^* - (37)$$

which:

$$A_x^* = \frac{A^{(0)}}{\sqrt{2}} e^{-i(\omega t - kz)} - (38)$$

$$A_y^* = -i \frac{A^{(0)}}{\sqrt{2}} e^{-i(\omega t - kz)} - (39)$$

From eq. (38):

$$\frac{\partial A_x^*}{\partial z} = i k A_x^* - (40)$$

From eq. (39)

$$\frac{\partial A_y^*}{\partial z} = k A_x^* - (41)$$

so

$$\omega_z = i k = k - (42)$$

The only possible solution is

$$\omega_z = 0 - (43)$$

which is eq. (31) Q.E.D.

hence the system rigorously obeys the conservation of antisymmetry.

b) The vector potential is:

$$\underline{A}^* = \underline{A}^{(2)} = \frac{\underline{A}^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} \quad (44)$$

and the spin connection vector is:

$$\underline{\omega} = \underline{\omega}^{(1)} = \frac{\kappa}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i\phi} \quad (45)$$

Eq. (24) is named the spin cyclic theorem.
 The other antisymmetry equations are:

$$\underline{E} = -\underline{\nabla} \phi + \underline{\omega} \phi = -\frac{\partial \underline{A}}{\partial t} - \omega_0 \underline{A} \quad (46)$$

and

$$\frac{1}{c^2} \left(\frac{\partial}{\partial t} + \omega_0 \right) \phi = (\underline{\nabla} - \underline{\omega}) \cdot \underline{A} \quad (47)$$

For the longitudinal components:

$$\underline{E}^{(3)} = \underline{\omega}^{(3)} = \underline{A}^{(3)} = 0 \quad (48)$$

so it follows:

$$\underline{\nabla} \phi = 0, \quad (49)$$

and

$$\left(\frac{\partial}{\partial t} + \omega_0 \right) \phi = 0, \quad (50)$$

a possible solution of which is:

$$\phi^{(3)} = 0 \quad (51)$$

For the transverse components:

$$7) \quad \underline{A}^* = \underline{A}^{(2)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} - (52)$$

So:

$$\underline{E}^{(2)} = - \frac{\partial \underline{A}^{(2)}}{\partial t} - \omega_0 \underline{A}^{(2)} - (53)$$

$$= - \frac{i\omega A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} - \omega_0 \underline{A}^{(2)}$$

$$= - \frac{i\omega A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} - \frac{\omega_0 A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi}$$

$$= \frac{\omega A^{(0)}}{\sqrt{2}} (-i\underline{i} + \underline{j}) e^{-i\phi} - \frac{\omega_0 A^{(0)}}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi}$$

In general, ω_0 may be found from:

$$\underline{B}^{(2)} = \underline{\nabla} \times \underline{A}^{(2)} - (54)$$

and

$$\underline{\nabla} \times \underline{E}^{(2)} + \frac{\partial \underline{B}^{(2)}}{\partial t} = \underline{0} - (55)$$

This is best done by computer algebra.
 The Lorenzian constraint for $\underline{A}^{(2)}$ is:

$$\frac{1}{c^2} \left(\frac{\partial}{\partial t} + \omega_0 \right) \phi = (\underline{\nabla} - \underline{\omega}) \cdot \underline{A}^* - (56)$$

which:

$$\underline{\nabla} \cdot \underline{A}^* = 0 - (57)$$

and:

$$\underline{\omega} \cdot \underline{A}^* = \frac{\omega^{(0)} A^{(0)}}{2} (\underline{i} - i\underline{j}) \cdot (\underline{i} + i\underline{j})$$

$$= \omega^{(0)} A^{(0)} \quad - (58)$$

So

$$\frac{1}{c^2} \left(\frac{\partial}{\partial t} + \omega_0 \right) \phi = \omega^{(0)} A^{(0)} = \kappa A^{(0)} \quad - (59)$$

and ϕ may be found.

Another procedure is to evaluate ϕ^* from:

$$\underline{\nabla} \cdot \underline{E}^* = \square \phi^* = 0 \quad - (60)$$

for the transverse plane write \underline{E}^* . So:

$$\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \right) \phi^* = 0 \quad - (61)$$

A solution is:

$$\phi^* = \phi^{(0)} \exp(-i(\omega t - \kappa z))$$

$$\kappa = \omega c \quad - (63)$$

Find ω_0 from

$$\frac{1}{c^2} \left(\frac{\partial}{\partial t} + \omega_0 \right) \phi = \kappa A^{(0)} \quad - (64)$$

Now adopt the procedure of Note 388(6),

and define:

$$\underline{E} = -\underline{\nabla} \phi + \underline{\omega} \phi := -\underline{\nabla} \phi - \frac{\partial \underline{A}}{\partial t} (\text{total}) \quad - (65)$$

So:
$$-\frac{\partial A(\text{total})}{\partial t} = \underline{\omega} \phi - (66)$$

where
$$\underline{A}(\text{total}) = \underline{A}^* + \underline{A}_1^* - (67)$$

and:
$$\underline{\nabla} \times \underline{A}_1^* = -\underline{\omega} \times \underline{A}^* - (68)$$

so
$$\underline{B}^* = \underline{\nabla} \times \underline{A}(\text{total}) - (69)$$

From eqs. (65) and (69):
$$\underline{\nabla} \times \underline{E}^* + \frac{\partial \underline{B}^*}{\partial t} = \underline{0} - (70)$$

From eq. (66):
$$\underline{A}^*(\text{total}) = -\int \underline{\omega} \phi^* dt + \underline{A}_2^* - (71)$$

so
$$\underline{A}_2^* = \underline{A}(\text{total}) + \int \underline{\omega} \phi^* dt - (72)$$

This procedure effectively solves eqs. (46) and (47) simultaneously so all the antisymmetry laws are satisfied.
