

319(3): Development of Note 319(2)

The force in EFE2 is defined by:

$$\underline{F} = m\underline{g} = -\underline{\nabla} U - \frac{\partial \underline{p}}{\partial t} - 2\underline{U}\underline{\omega} + 2c\omega_0 \underline{p} \quad (1)$$

with antisymmetry law:

$$-\underline{\nabla} U - 2\underline{U}\underline{\omega} = -\frac{\partial \underline{p}}{\partial t} + 2c\omega_0 \underline{p} \quad (2)$$

-(3)

so:

$$\underline{F} = 2 \left( -\underline{\nabla} U - 2\underline{U}\underline{\omega} \right) = 2 \left( -\frac{\partial \underline{p}}{\partial t} + 2c\omega_0 \underline{p} \right)$$

The Newtonian theory is classical, so it is not relativistic and there is no spin connection. So the Newtonian theory can be thought of as a limit

$$\underline{\nabla} U = \frac{\partial \underline{p}}{\partial t} \quad (4)$$

Eq. (4) is the Newtonian equivalence principle:

$$\underline{F} = m\underline{g} = -\frac{mM G}{r^2} \underline{e}_r \quad (5)$$
$$= -\frac{\partial \underline{p}}{\partial t}$$

Possible solutions of eq. (3) are:

$$\frac{\partial \underline{p}}{\partial t} = 2\underline{U}\underline{\omega} \quad (6)$$

and

$$\underline{\nabla} U = -2c\omega_0 \underline{p} \quad (7)$$

With assumption (6) and (7):

$$\begin{aligned} \underline{F} &= 2 \left( -\underline{\nabla} U - \frac{\partial \underline{p}}{\partial t} \right) = 2 \left( -2\underline{U}\underline{\omega} + 2c\omega_0 \underline{p} \right) \\ &= m \underline{g} \quad (8) \end{aligned}$$

This is eq. (7) of Note 315(2), Q.E.D.

If it is assumed in eq. (3) that:

$$\underline{\nabla} U = 2\underline{U}\underline{\omega} \quad (9)$$

and

$$-\frac{\partial \underline{p}}{\partial t} = 2c\omega_0 \underline{p} \quad (10)$$

it follows that:

$$\underline{\nabla} U = \frac{\partial \underline{p}}{\partial t} \quad (11)$$

so a Newtonian-like equivalence principle can be derived for conditions (9) and (10). Eq. (11) is however an equivalence principle of generally covariant unified field theory under the assumptions (9) and (10).

2) It follows that the most general equivalence principle is the antisymmetry equation (3). This, further, leads to a generalization of the other equivalence principle of physics.

The solutions of eqns. (9) and (10) can be written as:

$$\underline{\nabla} = 2\underline{\omega} \quad - (12)$$

and

$$\frac{1}{c} \frac{d}{dt} = -2\omega_0 \quad - (13)$$

i.e.

$$d_\mu = \left( \frac{1}{c} \frac{d}{dt}, \underline{\nabla} \right) \quad - (14)$$

$$= -2\omega_\mu = -2 \left( \omega_0, -\underline{\omega} \right)$$

Therefore under the condition:

$$\boxed{\omega_\mu = -\frac{1}{2} d_\mu} \quad - (15)$$

The most general equivalence principle reduces to eq. (11), i.e. to a Newtonian-like result but part of a generally covariant unified field theory. Conversely this indicates that the Newtonian equivalence principle is a manifestation of the more general result (3) under the condition (15).

4) The general equivalence principle (3) means that gravitational forces vanish when:

$$-\underline{\nabla} U - 2\underline{U} \underline{\omega} = - \frac{d\underline{p}}{dt} + 2c\underline{\omega} \cdot \underline{p} = \underline{0} \quad -(16)$$

Eq. (16) means:

$$\underline{\nabla} U = -2\underline{U} \underline{\omega} \quad -(17)$$

and

$$\frac{d\underline{p}}{dt} = 2c\underline{\omega} \cdot \underline{p} \quad -(18)$$

i.e.

$$\underline{\omega}_\mu = \frac{1}{2} \dot{\underline{p}}_\mu \quad -(19)$$

This is the opposite condition to eq. (15)

The basic gravitational field equations of GECF2 are:

$$\underline{\nabla} \cdot \underline{\Omega} = 0 \quad -(16)$$

$$\underline{\nabla} \times \underline{g} + \frac{d\underline{\Omega}}{dt} = \underline{0} \quad -(17)$$

$$\underline{\nabla} \cdot \underline{g} = \underline{\kappa} \cdot \underline{g} = 4\pi G \rho_m \quad -(18)$$

$$\underline{\nabla} \times \underline{\Omega} - \frac{1}{c^2} \frac{d\underline{g}}{dt} = \frac{4\pi G}{c^2} \underline{J}_m = \underline{\kappa} \times \underline{\Omega} \quad -(19)$$

These field equations have a Lorentz-like

5) covariant but are equivalent of a generally covariant unified field theory (ECE2)

This means that they can be written in a Minkowski-like metric, always bearing in mind that they are equivalent to a general metric.

This realization links ECE2 to  $\alpha$  theory, and is a link to papers such as UFT 216 and UFT 261, and the development of orbital theory with a Minkowski-like metric. The velocity from a Minkowski-like metric is

$$v^2 = \left(\frac{dr}{dt}\right)^2 + r^2 \left(\frac{d\theta}{dt}\right)^2 \quad (20)$$

for a planar orbit. For an observed orbit of type:

$$r = \frac{d}{1 + \epsilon \cos(\alpha\theta)} \quad (21)$$

of  $\alpha$  theory the deflection of light due to gravitation for eqs. (20) to (21) is obtained from the velocity:

$$v^2 = \frac{L^2}{R_0 m^2} \left[ \frac{\alpha^2}{d} (1 + \epsilon) - \frac{1}{R_0} (\alpha^2 - 1) \right] \quad (22)$$

In the Newtonian limit:

$$x = 1 \quad (23)$$

So

$$v^2 = \frac{L^2 (1+\epsilon)}{R_0 m^2 d} \quad (24)$$

where  $R_0$  is the distance of closest approach.

The limit (24) should be regarded as the limit of a generally covariant unified field theory when:

$$r = \frac{d}{1+\epsilon \cos \theta} \quad (25)$$

which is the conical section of Newtonian theory. Light deflected by the sun is described by the hyperbola:

$$\epsilon > 1. \quad (26)$$

In this limit,

$$v^2 = \frac{M G}{R_0^2} \quad (27)$$

So

$$v^2 = \frac{M G}{R_0^2} \quad (28)$$

where

$$R_0 = \frac{d}{1+\epsilon} \quad (29)$$

In this limit the angle of deflection is:

$$2\phi = \frac{2}{\epsilon} = 2 \left( \frac{R_0 v^2}{M G} - 1 \right)^{-1}$$

$$\sim \frac{2 M G}{R_0 v^2} \quad (30)$$

1) The experimentally observed result is:

$$2\phi = \frac{4mG}{R_0 c^2} \quad - (31)$$

This result can be expressed as:

$$2\phi = - \frac{\Phi}{c^2} \quad - (31)$$

where

$$\Phi = - \frac{4mG}{R_0} \quad - (32)$$

Eq. (30) is obtained as the limit of a more general expression: - (33)

$$2\phi = \frac{2}{c} \left[ \frac{n^2 d R_0}{x c^2 L^2} \left( v^2 + \frac{L^2}{n^2} \left( \frac{1-x^2}{R_0^2} \right) \right) - 1 \right]$$

corresponding to  $r = \frac{d}{1 + \epsilon \cos(x\theta)} \quad - (34)$

so the experimental result is obtained from:

$$\left( \frac{v^2}{c^2} + \frac{L^2}{n^2 R_0^2} \left( \frac{1-x^2}{x^2} \right) \right) = \frac{2}{c^2} \quad - (35)$$

This equation defines a Lorentz factor of

o) the Lorentz-like eqs (16) to (19).

So

$$v^2 + v_1^2 = \frac{\alpha c^2}{2} \quad - (36)$$

where

$$v_1^2 = \frac{L^2 (1 - \alpha^2)}{n^2 R_0^2} \quad - (37)$$

To an excellent approximation:

$$\alpha \approx 1 \quad - (38)$$

So

$$\boxed{c^2 = 2v^2} \quad - (39)$$

This is the condition needed for:

$$2\phi = \frac{4MG}{R_0 c^2} = \frac{2MG}{R_0 v^2} \quad - (40)$$

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