

$$(a) - \theta = \frac{1}{2} \left(\frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial y} \right) \right)$$

$$(b) - \theta = \frac{1}{2} \left(\frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial y} \right) \right)$$

$$(c) - \theta = \frac{1}{2} \left(\frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial x} \right) \right)$$

$$(d) - \theta = \frac{1}{2} \left(\frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial y} \right) \right)$$

(e) -

$$\theta = \frac{1}{2} \left(\frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial y} \right) \right)$$

(f) - $\{ \epsilon \} | \theta = 0$

: no singularities for θ if $\{ \epsilon \} \neq 0$ (T)

$$\theta = \frac{1}{2} \left(\frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial x} \right) \right)$$

(g) - : (f) from (e)

$$\frac{\partial u}{\partial x} = \frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} + \frac{\partial v}{\partial x} \right) \quad \text{: o}$$

(e) - This is summarized as perfect θ in Eq.

$$(e) - \frac{\partial u}{\partial x} = \frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \quad \text{: else}$$

$$(f) - \theta = \frac{\partial}{\partial t} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right)$$

: perfect θ if $\{ \epsilon \} \neq 0$

Exercise 15.3: Develop a numerical method for the field equations from

15.3: Develop a numerical method for the field equations from

• 2. If you grant others such rights (S1) or (e1). It is also subject to
• (II) - if anyone goes against such rights (S1) or (e1). It is also subject to

$$0 = \sqrt{m_0} \frac{1}{\sqrt{q}} T + \sqrt{m_0} \frac{1}{\sqrt{q}} T + \sqrt{m_1} \frac{1}{\sqrt{q}} T + \sqrt{m_0} \frac{1}{\sqrt{q}} T$$

(71) -

$$0 = \sum_{n=0}^{\infty} \frac{1}{q} e^{nw} - \sum_{n=1}^{\infty} \frac{1}{q} e^{nw} + \sum_{n=1}^{\infty} \frac{1}{q} e^{nw} + \dots$$

(81)

$$0 = \sqrt{q} \frac{1}{\sqrt{q}} + \sqrt{q} \frac{1}{\sqrt{q}} + \sqrt{q} \frac{1}{\sqrt{q}} + \sqrt{q} \frac{1}{\sqrt{q}}$$

$$0 = \operatorname{erf} \frac{1}{\sqrt{\pi}} \frac{1}{q} + \operatorname{erf} \frac{1}{\sqrt{\pi}} \frac{1}{q} + \operatorname{erf} \frac{1}{\sqrt{\pi}} \frac{1}{q} + \operatorname{erf} \frac{1}{\sqrt{\pi}} \frac{1}{q}$$

If it is difficult to find a suitable
substitution, one may add the following
substances and substituents. The following
substituents are substituted in the following
order: (i) by nitrogen, (ii) by halogen, (iii)

$$(11) - o =$$

$$\left(\sqrt[n]{\frac{1}{q}} + \sqrt[n]{\frac{1}{q}} + \sqrt[n]{\frac{1}{q}} + \dots + \sqrt[n]{\frac{1}{q}} \right) +$$

$$\left(\text{exp} \frac{1}{q} + \text{exp} \frac{1}{q} + \text{exp} \frac{1}{q} + \text{exp} \frac{1}{q} \right) e^{\lambda} +$$

$$\left(\arcsin \frac{1}{q} + \arcsin \frac{1}{q} + \arcsin \frac{1}{q} + \arcsin \frac{1}{q} \right) \alpha +$$

$$\left(\sin \frac{1}{q} + \cos \frac{1}{q} + \tan \frac{1}{q} + \cot \frac{1}{q} \right) \approx 0$$

Adds $\alpha \wedge \beta$ to (α)

$$\begin{aligned}
 & T_{\mu_0}^b \tilde{T}^{0\mu_0} + T_{\mu_0}^b \tilde{T}^{1\mu_1} + T_{\mu_0}^b \tilde{T}^{2\mu_2} + T_{\mu_0}^b \tilde{T}^{3\mu_3} \\
 & + T_{\mu_1}^b \tilde{T}^{1\mu_0} + T_{\mu_1}^b \tilde{T}^{1\mu_1} + T_{\mu_1}^b \tilde{T}^{1\mu_2} + T_{\mu_1}^b \tilde{T}^{1\mu_3} \\
 & + T_{\mu_2}^b \tilde{T}^{2\mu_0} + T_{\mu_2}^b \tilde{T}^{2\mu_1} + T_{\mu_2}^b \tilde{T}^{2\mu_2} + T_{\mu_2}^b \tilde{T}^{2\mu_3} \\
 & + T_{\mu_3}^b \tilde{T}^{3\mu_0} + T_{\mu_3}^b \tilde{T}^{3\mu_1} + T_{\mu_3}^b \tilde{T}^{3\mu_2} + T_{\mu_3}^b \tilde{T}^{3\mu_3} - (16) \\
 & = 0
 \end{aligned}$$

In UFT 314 it was assumed that the sum of the diagonals in Eq. (16) is zero:

$$\begin{aligned}
 & T_{\mu_0}^b \tilde{T}^{a\mu_0} + T_{\mu_1}^b \tilde{T}^{a\mu_1} + T_{\mu_2}^b \tilde{T}^{a\mu_2} + T_{\mu_3}^b \tilde{T}^{a\mu_3} \\
 & = 0 - (17)
 \end{aligned}$$

For electrodynamics and in vector notation this lead to the result:

$$\underline{E}^{(1)} \cdot \underline{B}^{(2)} + \underline{E}^{(2)} \cdot \underline{B}^{(1)} = 0. - (18)$$

Eq. (18) was verified for plane waves.

The assumption leading to Eq. (17) means that:

$$\begin{aligned}
 & T_{\mu_0}^b (\tilde{T}^{a\mu_1} + \tilde{T}^{a\mu_2} + \tilde{T}^{a\mu_3}) \\
 & + T_{\mu_1}^b (\tilde{T}^{a\mu_0} + \tilde{T}^{a\mu_2} + \tilde{T}^{a\mu_3}) \\
 & + T_{\mu_2}^b (\tilde{T}^{a\mu_0} + \tilde{T}^{a\mu_1} + \tilde{T}^{a\mu_3}) \\
 & + T_{\mu_3}^b (\tilde{T}^{a\mu_0} + \tilde{T}^{a\mu_1} + \tilde{T}^{a\mu_2}) = 0 - (19)
 \end{aligned}$$

4) Eq. (19) is implied by eqs. (16) and (17) and can be used to derive more field relations in electrodynamics and gravitation. In electrodynamics:

$$F_{\mu\nu}^b (\tilde{F}^{a\mu\nu} + \tilde{F}^{a\mu\nu 2} + \tilde{F}^{a\mu\nu 3}) + F_{\mu 1}^b (\tilde{F}^{a\mu\nu} + \tilde{F}^{a\mu\nu 2} + \tilde{F}^{a\mu\nu 3}) = 0 \quad -(20)$$

$$+ F_{\mu 2}^b (\tilde{F}^{a\mu\nu} + \tilde{F}^{a\mu\nu 1} + \tilde{F}^{a\mu\nu 3})$$

$$+ F_{\mu 3}^b (\tilde{F}^{a\mu\nu} + \tilde{F}^{a\mu\nu 1} + \tilde{F}^{a\mu\nu 2})$$

where:

$$F_{\mu\nu}^b = \begin{bmatrix} 0 & \tilde{E}_x^b & \tilde{E}_y^b & \tilde{E}_z^b \\ -\tilde{E}_x^b & 0 & -cB_z^b & cB_y^b \\ -\tilde{E}_y^b & cB_z^b & 0 & -cB_x^b \\ -\tilde{E}_z^b & -cB_y^b & cB_x^b & 0 \end{bmatrix} \quad -(21)$$

$$\tilde{F}^{a\mu\nu} = \begin{bmatrix} 0 & -cB_x^a & -cB_y^a & -cB_z^a \\ cB_x^a & 0 & \tilde{E}_z^a & -\tilde{E}_y^a \\ cB_y^a & -\tilde{E}_z^a & 0 & \tilde{E}_x^a \\ cB_z^a & \tilde{E}_y^a & -\tilde{E}_x^a & 0 \end{bmatrix} \quad -(22)$$

Eq. (20) is evaluated by summing over repeated μ indices using eqs. (21) and (22). It is proven as follows that eq. (20) is true given eqs. (21) and (22).

5) Proof

Eq. (20) is:

$$\begin{aligned} & F_{10}^b (\tilde{F}^{a12} + \tilde{F}^{a13}) + F_{20}^b (\tilde{F}^{a21} + \tilde{F}^{a23}) + F_{30}^b (\tilde{F}^{a31} + \tilde{F}^{a32}) \\ & + F_{01}^b (\tilde{F}^{a02} + \tilde{F}^{a03}) + F_{21}^b (\tilde{F}^{a20} + \tilde{F}^{a23}) + F_{31}^b (\tilde{F}^{a30} + \tilde{F}^{a32}) \\ & + F_{02}^b (\tilde{F}^{a01} + \tilde{F}^{a03}) + F_{12}^b (\tilde{F}^{a10} + \tilde{F}^{a13}) + F_{32}^b (\tilde{F}^{a30} + \tilde{F}^{a31}) \\ & + F_{03}^b (\tilde{F}^{a01} + \tilde{F}^{a02}) + F_{13}^b (\tilde{F}^{a10} + \tilde{F}^{a12}) + F_{23}^b (\tilde{F}^{a20} + \tilde{F}^{a21}) \\ & = 0 \quad - (23) \end{aligned}$$

Using eqs. (21) and (22), eq. (23) is:

$$\begin{aligned} & -E_x^b (E_z^a - E_y^a) - E_y^b (-E_z^a + E_x^a) - E_z^b (E_y^a - E_x^a) \\ & + E_x^b (-cB_y^a - cB_z^a) + cB_z^b (cB_y^a + E_x^a) - cB_y^b (cB_z^a - E_x^a) \\ & - E_y^b (cB_x^a + cB_z^a) - cB_z^b (cB_x^a - E_y^a) + cB_x^b (cB_z^a + E_y^a) \\ & + E_z^b (-cB_x^a - cB_y^a) + cB_y^b (cB_x^a + E_z^a) - cB_x^b (cB_y^a - E_z^a) \\ & = 0 \quad - (24) \end{aligned}$$

Q.E.D., if $a = b$.

Therefore eqs. (17) and (18) are also true, Q.E.D.

The following four ident. k's are also true:

$$6) \quad F_{\mu_0}^b (\tilde{F}^{a\mu_1} + \tilde{F}^{a\mu_2} + \tilde{F}^{a\mu_3}) = 0 \quad - (25)$$

$$F_{\mu_1}^b (\tilde{F}^{a\mu_0} + \tilde{F}^{a\mu_2} + \tilde{F}^{a\mu_3}) = 0 \quad - (26)$$

$$F_{\mu_2}^b (\tilde{F}^{a\mu_0} + \tilde{F}^{a\mu_1} + \tilde{F}^{a\mu_3}) = 0 \quad - (27)$$

$$F_{\mu_3}^b (\tilde{F}^{a\mu_0} + \tilde{F}^{a\mu_1} + \tilde{F}^{a\mu_2}) = 0 \quad - (28)$$

With summation over the μ indices in each case.

Eqs. (25) to (28) are true if:

$$a = b \quad - (29)$$

In this case, eq. (18) reduces for each sense of polarization to:

$$\boxed{\underline{E} \cdot \underline{B} = 0} \quad - (30)$$

This is the well known result that electric and magnetic fields are perpendicular in free space.

For example, if:

$$\underline{E} = \frac{\underline{E}^{(0)}}{\sqrt{2}} (i - j) e^{i\phi} \quad - (31)$$

then it follows from:

$$\nabla \times \underline{E} + \frac{\partial \underline{B}}{\partial t} = \underline{0} \quad - (32)$$

7) That: $\underline{B} = \frac{\underline{B}^{(0)}}{\sqrt{2}} (i\underline{i} + j) e^{i\phi} - (33)$

From eqs. (31) and (33) eq. (30) follows, Q.E.D.

Note that if:

$$\underline{E}^{(1)} = \frac{\underline{E}^{(0)}}{\sqrt{2}} (i - i\underline{j}) e^{i\phi} - (34)$$

then

$$\underline{B}^{(2)} = \frac{\underline{B}^{(0)}}{\sqrt{2}} (-i\underline{i} + j) e^{-i\phi} - (35)$$

Also:

$$\underline{E}^{(2)} = \frac{\underline{E}^{(0)}}{\sqrt{2}} (i + i\underline{j}) e^{-i\phi} - (36)$$

and

$$\underline{B}^{(1)} = \frac{\underline{B}^{(0)}}{\sqrt{2}} (i\underline{i} + j) e^{i\phi} - (37)$$

So

$$\underline{E}^{(1)} \cdot \underline{B}^{(2)} + \underline{E}^{(2)} \cdot \underline{B}^{(1)} = 0 - (38)$$

Therefore

$$(\underline{E} \cdot \underline{B})^a = (\underline{E} \cdot \underline{B})^b = 0 - (39)$$

and

$$\underline{E}^a \cdot \underline{B}^b + \underline{E}^b \cdot \underline{B}^a = 0 - (40)$$