

275(1) : Simplest Formulation of the 3D Orbit
for the Inverse Square Law.

The two basic equations of the orbit are:

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1 \quad - (1)$$

and

$$z^2 = \left(1 - \frac{L_z^2}{L^2}\right) y^2 \quad - (2)$$

here:

$$x = ae + r \cos \beta \quad - (3)$$

$$y = r \sin \beta \quad - (4)$$

$$r = \frac{d}{1 + e \cos \beta} \quad - (5)$$

$$d = a(1 - e^2) \quad - (6)$$

$$e^2 = 1 - \frac{b^2}{a^2} \quad - (7)$$

$$H = \frac{1}{2} m (\dot{r}^2 + \dot{\beta}^2 r^2) - \frac{k}{r} \quad - (8)$$

$$L = \frac{1}{2} m (\dot{r}^2 + \dot{\beta}^2 r^2) + \frac{k}{r} \quad - (9)$$

$$k = mM\epsilon \quad - (10)$$

$$\dot{\beta}^2 = \dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2 \quad - (11)$$

The angular equations are:

2)

$$\cos^2 \theta = \left(\frac{L^2 - L_z^2}{L^2} \right) \sin^2 \beta \quad - (12)$$

$$= \left(1 - \left(\frac{L_z}{L} \right)^2 \right) \sin^2 \beta$$

and

$$\tan \phi = \left(\frac{L_z}{L} \right) \tan \beta \quad - (13)$$

i.e.

$$\cos^2 \beta = \frac{\cos^2 \phi}{\cos^2 \phi + \left(\frac{L}{L_z} \right)^2 \sin^2 \phi} \quad - (14)$$

The angular relations (12) to (14) depend only on eq. (11), and are true for any potential energy U of the Hamiltonian:

$$H = T + U \quad - (15)$$

and Lagrangian: $L = T - U. \quad - (16)$

Divide eq. (2) by a distance c yet to be defined:

$$\frac{z^2}{c^2} = \left(1 - \frac{L_z}{L} \right) \frac{y^2}{c^2} \quad - (17)$$

3) Add equations (1) and (17):

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} = 1 + \left(1 - \frac{L_2}{L}\right) \frac{y^2}{c^2} \quad - (18)$$

This is the ellipsoidal orbit:

$$\frac{x^2}{a^2} + y^2 \left(\frac{1}{b^2} - \frac{1}{c^2} \left(1 - \frac{L_2}{L}\right) \right) + \frac{z^2}{c^2} = 1 \quad - (19)$$

observed in three dimensional galaxies.

Reduction to Spherical Orbit.

The spherical orbit is:

$$x^2 + y^2 + z^2 = r^2 \quad - (20)$$

and eq. (19) becomes eq. (20) when:

$$\frac{1}{a^2} = \frac{1}{c^2} = \frac{1}{b^2} - \frac{1}{c^2} \left(1 - \frac{L_2}{L}\right) = \frac{1}{r^2} \quad - (21)$$

i.e.

$$a = c = r \quad - (22)$$

$$\frac{1}{b^2} - \frac{1}{c^2} \left(1 - \frac{L_2}{L}\right) = \frac{1}{r^2} \quad - (23)$$

4) From eq. (22):

$$c = r \quad - (24)$$

so:

$$\frac{1}{b^2} - \frac{1}{r^2} \left(1 - \frac{L_2}{L} \right) = \frac{1}{r^2} \quad - (25)$$

and

$$\begin{aligned} \frac{1}{b^2} &= \frac{1}{r^2} \left(1 + 1 - \frac{L_2}{L} \right) \\ &= \frac{1}{r^2} \left(2 - \frac{L_2}{L} \right) \quad - (26) \end{aligned}$$

Therefore the spherical orbit (20) is obtained

when:

$$a = c = r \quad - (27)$$

and

$$b = r \left(2 - \frac{L_2}{L} \right)^{-1/2} \quad - (28)$$

Elongated Ellipsoidal Orbit

In this case:

$$a^2 \gg b^2 \sim c^2 \quad - (29)$$

The spherical and elongated ellipsoidal orbits are observed experimentally in three dimensional galaxies.

5) The semi major axis a and semi minor axis b of the seta ellipse are defined by:

$$d = \frac{L^2}{mk} \quad - (30)$$

and

$$e^2 = 1 + \frac{2EL^2}{mk^2} \quad - (31)$$

where L is the total angular momentum and E the total energy, two constants of motion. So:

$$d = \frac{L^2}{mk} = a(1-e^2) \quad - (32)$$

and

$$e^2 = 1 - \frac{b^2}{a^2} = 1 + \frac{2EL^2}{mk^2} \quad - (33)$$

It follows that:

$$E = -\frac{b^2}{a^2} \frac{mk^2}{2L^2} \quad - (34)$$

and

$$d = a(1-e^2) = \frac{b^2}{a} \quad - (35)$$

So:

$$L^2 = mkd = mk \frac{b^2}{a} \quad - (36)$$

From eqs (34) and (36):

$$E = -\frac{k}{2a} \quad - (37)$$

Therefore for β let a ellipse:

$$0 < E < 1 \quad - (38)$$

and

$$V_{\min} < E < 0 \quad - (39)$$

where

$$V = -\frac{k}{r} + \frac{L^2}{2mr^2} \quad - (40)$$

Therefore:

$$|E| = \frac{k}{2a} \quad - (41)$$

and

$$a = \frac{k}{2|E|} = \frac{mMG}{2|E|} \quad - (42)$$

It follows that:

$$1 - e^2 = \frac{2|E|d}{k} \quad - (43)$$

and that:

$$b = \frac{d}{(1-e^2)^{1/2}} = \frac{L}{(2m|E|)^{1/2}} \quad - (44)$$