

258(5) : Proof of the Free Space Beltrami Equations
in ECE Theory.

Start with the Gauss law:

$$\underline{\nabla} \cdot \underline{B}^a = 0 \quad - (1)$$

and use the identity: $\underline{\nabla} \cdot \underline{\nabla} \times \underline{B}^a = 0 \quad - (2)$

It follows that: $\underline{\nabla} \times \underline{B}^a = \kappa \underline{B}^a \quad - (3)$

and $\underline{B}^a = \frac{1}{\kappa} \underline{\nabla} \times \underline{B}^a \quad - (4)$

Therefore:
$$\begin{aligned} \underline{\nabla} \cdot \underline{\nabla} \times \underline{B}^a &= \underline{\nabla} \cdot (\kappa \underline{B}^a) \\ &= \kappa \underline{\nabla} \cdot \underline{B}^a + \underline{B}^a \cdot \underline{\nabla} \kappa \\ &= 0 \end{aligned} \quad - (5)$$

Therefore: $\underline{B}^a \cdot \underline{\nabla} \kappa = 0 \quad - (6)$

Also: $\underline{\nabla} \cdot \underline{B}^a = \underline{\nabla} \cdot \left(\frac{1}{\kappa} \underline{\nabla} \times \underline{B}^a \right) = 0 \quad - (7)$

so:
$$\begin{aligned} \underline{\nabla} \times \underline{B}^a \cdot \underline{\nabla} \left(\frac{1}{\kappa} \right) &= 0 \\ &= \kappa \underline{B}^a \cdot \underline{\nabla} \left(\frac{1}{\kappa} \right) \end{aligned} \quad - (8)$$

From eqs. (6) and (8):

$$\boxed{\frac{1}{\kappa} \underline{\nabla} \kappa = \underline{\nabla} \left(\frac{1}{\kappa} \right)} \quad - (9)$$

2) $\underline{\epsilon}$: $\kappa = \kappa(x, y, z) - (10)$

it is governed by eq. (9). If κ is a constant then eq. (9) is solved automatically. Usually in Beltrami theory, κ is a constant independent of x, y and z .

In free space:

$$\underline{\nabla} \cdot \underline{E}^a = 0 - (11)$$

so it follows immediately that:

$$\boxed{\underline{\nabla} \times \underline{E}^a = \kappa \underline{E}^a} - (12)$$

Also in free space:

$$\underline{\nabla} \times \underline{E}^a + \frac{\partial \underline{B}^a}{\partial t} = \underline{0} - (13)$$

$$\underline{\nabla} \times \underline{B}^a - \frac{1}{c^2} \frac{\partial \underline{E}^a}{\partial t} = \underline{0} - (14)$$

so:

$$\frac{\partial \underline{B}^a}{\partial t} = -\kappa \underline{E}^a - (15)$$

$$\frac{1}{c^2} \frac{\partial \underline{E}^a}{\partial t} = \kappa \underline{B}^a - (16)$$

It follows that:

$$\frac{\partial^2 \underline{E}^a}{\partial t^2} = -\omega^2 \underline{E}^a - (17)$$

3) and

$$\frac{\partial^2 \underline{B}^a}{\partial t^2} = -\omega^2 \underline{B}^a \quad (18)$$

so

$$\underline{E}^a = \underline{E}_0^a e^{i\omega t} \quad (19)$$

$$\underline{B}^a = \underline{B}_0^a e^{i\omega t} \quad (20)$$

in free space, where:

$$\underline{\nabla} \times \underline{E}^a = \kappa \underline{E}^a \quad (21)$$

$$\underline{\nabla} \times \underline{B}^a = \kappa \underline{B}^a \quad (22)$$

Note carefully that eqs (19) - (22) are true for all electromagnetic fields in free space.

From the Coulomb identity:

$$\underline{\nabla} \cdot (\underline{\omega}^a_b \times \underline{A}^b) = 0 \quad (23)$$

in free space. It follows that:

$$\underline{\nabla} \times (\underline{\omega}^a_b \times \underline{A}^b) = \kappa \underline{\omega}^a_b \times \underline{A}^b \quad (24)$$

The solutions of $\underline{\omega}^a_b \times \underline{A}^b$ are solutions of a Helmholtz equation in the absence of a magnetic monopole.

Eq. (23) implies:

$$4) \quad \underline{\omega}^a{}_b \cdot \underline{\nabla} \times \underline{A}^b = \underline{A}^b \cdot \underline{\nabla} \times \underline{\omega}^a{}_b - (25)$$

The magnetic field is defined by:

$$\underline{B}^a = \underline{\nabla} \times \underline{A}^a - \underline{\omega}^a{}_b \times \underline{A}^b - (26)$$

From eqns. (3) and (26):

$$\begin{aligned} \underline{\nabla} \times \underline{B}^a &= \underline{\nabla} \times (\underline{\nabla} \times \underline{A}^a - \underline{\omega}^a{}_b \times \underline{A}^b) - (27) \\ &= \kappa \underline{B}^a \\ &= \kappa (\underline{\nabla} \times \underline{A}^a - \underline{\omega}^a{}_b \times \underline{A}^b) \end{aligned}$$

$$\text{So:} \quad \underline{\nabla} \times (\underline{\nabla} \times \underline{A}^a) = \kappa \underline{\nabla} \times \underline{A}^a - (28)$$

$$\text{i.e.} \quad \boxed{\underline{\nabla} \times \underline{A}^a = \kappa \underline{A}^a} - (29)$$

From eqns. (25) and (29):

$$\boxed{\underline{\nabla} \times \underline{\omega}^a{}_b = \kappa \underline{\omega}^a{}_b} - (30)$$

Therefore in free space:

$$\underline{\nabla} \times \underline{B}^a = \kappa \underline{B}^a - (31)$$

$$\underline{\nabla} \times \underline{E}^a = \kappa \underline{E}^a - (32)$$

$$\underline{\nabla} \times \underline{A}^a = \kappa \underline{A}^a - (33)$$

$$\underline{\nabla} \times \underline{\omega}^a{}_b = \kappa \underline{\omega}^a{}_b - (34)$$

$$\underline{\nabla} \times (\underline{\omega}^a{}_b \times \underline{A}^b) = \underline{\omega}^a{}_b \times \underline{A}^b - (35)$$

5) These are general results for vacuum electro-magnetism of any kind. All these quantities have a rich variety of solutions, all with longitudinal components.

Plane Wave Solutions

In this case:

$$\underline{\omega}^a \underline{b} = \kappa \epsilon^{abc} \underline{A}^c / A^{(0)} - (36)$$

where:

$$\underline{A}^{(1)} \times \underline{A}^{(2)} = i A^{(0)} \underline{A}^{(3)*} - (37)$$

$$\underline{A}^{(3)} \times \underline{A}^{(1)} = i A^{(0)} \underline{A}^{(2)*} - (38)$$

$$\underline{A}^{(2)} \times \underline{A}^{(3)} = i A^{(0)} \underline{A}^{(1)*} - (39)$$

Here:

$$\underline{\nabla} \times \underline{A}^{(1)} = \kappa \underline{A}^{(1)} - (40)$$

$$\underline{\nabla} \times \underline{A}^{(2)} = \kappa \underline{A}^{(2)} - (41)$$

$$\underline{\nabla} \times \underline{A}^{(3)} = 0 \underline{A}^{(3)} - (42)$$

and

$$\underline{\nabla} \cdot \underline{A}^{(1)} = \underline{\nabla} \cdot \underline{A}^{(2)} = \underline{\nabla} \cdot \underline{A}^{(3)} = 0 - (43)$$

where

$$\underline{A}^{(1)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} - i \underline{j}) e^{i\phi} - (44)$$

$$\underline{A}^{(2)} = \frac{A^{(0)}}{\sqrt{2}} (\underline{i} + i \underline{j}) e^{-i\phi} - (45)$$

$$\underline{A}^{(3)} = A^{(0)} \underline{k} - (46)$$

b) The electromagnetic phase is:

$$\phi = \omega t - \kappa z - (47)$$

where ω is the angular frequency at t and κ the wave vector magnitude at z .

For plane waves the longitudinal component is a constant along z .

It is easily checked that eqs. (31) to (35) hold for plane waves. For example:

$$\underline{\nabla} \times (\underline{A}^{(2)} \times \underline{A}^{(3)}) = \kappa \underline{A}^{(2)} \times \underline{A}^{(3)} - (48)$$

From eq. (39), this equation is

$$\underline{\nabla} \times \underline{A}^{(2)} = \kappa \underline{A}^{(2)} - (49)$$

Self consistently:

$$\begin{aligned} \underline{\nabla} \times (\underline{A}^{(2)} \times \underline{A}^{(3)}) &= \underline{A}^{(2)} (\underline{\nabla} \cdot \underline{A}^{(3)}) - (\underline{\nabla} \cdot \underline{A}^{(2)}) \underline{A}^{(3)} \\ &+ (\underline{A}^{(2)} \cdot \underline{\nabla}) \underline{A}^{(3)} - (\underline{A}^{(3)} \cdot \underline{\nabla}) \underline{A}^{(2)} - (50) \\ &= - \underline{A}^{(2)} \frac{\partial \underline{A}^{(2)}}{\partial z} \\ &= \kappa \underline{A}^{(2)} \end{aligned}$$

Q.E.D.