

203(4) : Full Details of the Newtonian Calculation of the Perihelia Precession.

The basis of the calculation is:

$$\theta \rightarrow x\theta. \quad - (1)$$

The Lagrangian is:

$$\begin{aligned} \mathcal{L} &= \frac{1}{2} m \left(\left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d(x\theta)}{dt} \right)^2 \right) + \frac{mMG}{r} \\ &= \frac{1}{2} m \left(\left(\frac{dr}{dt} \right)^2 + x^2 r^2 \left(\frac{d\theta}{dt} \right)^2 \right) + \frac{mMG}{r} \quad - (2) \end{aligned}$$

where
$$U = -\frac{mMG}{r} \quad - (3)$$

The Euler Lagrange equation is:

$$\frac{\partial \mathcal{L}}{\partial \theta} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = 0 \quad - (4)$$

and the angular momentum is:

$$L = \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = x^2 m r^2 \frac{d\theta}{dt} \quad - (5)$$

The total energy is:

$$\begin{aligned} E &= \frac{1}{2} m \left(\left(\frac{dr}{dt} \right)^2 + x^2 r^2 \left(\frac{d\theta}{dt} \right)^2 \right) + U \\ &= \frac{1}{2} m \left(\frac{dr}{dt} \right)^2 + \frac{L^2}{2m x^2 r^2} + U \quad - (6) \end{aligned}$$

So:
$$\frac{dr}{dt} = \left(\frac{2}{m} \left(E - U - \frac{L^2}{2m x^2 r^2} \right) \right)^{1/2} \quad - (7)$$

Now we:

$$\frac{d\theta}{dr} = \frac{d\theta}{dt} \frac{dt}{dr} \quad - (8)$$

where
$$\frac{d\theta}{dt} = \frac{L}{m x^2 r^2} \quad - (9)$$

so:

$$\frac{dr}{d\theta} = \frac{m \omega^2 r^2}{L} \left(\frac{2}{m} \left(E - U - \frac{L^2}{2m \omega^2 r^2} \right) \right)^{1/2}$$

and $\left(\frac{dr}{d\theta} \right)^2 = \frac{2m \omega^4 r^4}{L^2} \left(E - U - \frac{L^2}{2m \omega^2 r^2} \right) \quad - (10)$

The astronomical observations give the precessing ellipse:

$$r = \frac{d}{1 + \epsilon \cos(\chi \theta)} \quad - (12)$$

so: $\frac{dr}{d\theta} = \frac{\chi \epsilon}{d} r^2 \sin(\chi \theta) \quad - (13)$

From eq. (12):

$$\cos(\chi \theta) = \frac{1}{\epsilon} \left(\frac{d}{r} - 1 \right) \quad - (14)$$

Use:

$$\sin^2(\chi \theta) + \cos^2(\chi \theta) = 1 \quad - (15)$$

so

$$\sin^2(\chi \theta) = 1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \quad - (16)$$

and $\left(\frac{dr}{d\theta} \right)^2 = \frac{\chi^2 \epsilon^2}{d^2} r^4 \left(1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \right) \quad - (17)$

Comparing eqs. (14) and (17):

$$\frac{2m \omega^4}{L^2} \left(E - U - \frac{L^2}{2m \omega^2 r^2} \right) = \frac{\chi^2 \epsilon^2}{d^2} \left(1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \right) \quad - (18)$$

So:

$$\begin{aligned} \frac{2mx^2}{L^2} \left(E - u - \frac{L^2}{2mx^2 r^2} \right) &= \frac{\epsilon^2}{d^2} \left(1 - \frac{1}{\epsilon^2} \left(\frac{d}{r} - 1 \right)^2 \right) \\ &= \frac{\epsilon^2}{d^2} \left(1 - \frac{1}{\epsilon^2 r^2} (d-r)^2 \right) \\ &= \frac{\epsilon^2}{d^2 r^2} \left(r^2 - \frac{1}{\epsilon^2} (d-r)^2 \right) \quad - (19) \end{aligned}$$

So:

$$\begin{aligned} \frac{\epsilon^2}{d^2} \left(r^2 - \frac{1}{\epsilon^2} (d-r)^2 \right) &= \frac{2mx^2 r^2}{L^2} \left(E - u - \frac{L^2}{2mx^2 r^2} \right) \\ &= \frac{\epsilon^2}{d^2} \left(r^2 - \frac{1}{\epsilon^2} (d^2 - 2dr + r^2) \right) \\ &= \frac{2mx^2 E}{L^2} r^2 + \frac{2mx^2}{L^2} m M G r - 1 \quad - (20) \end{aligned}$$

Comparing terms, eqn. (20) is:

$$= \frac{\epsilon^2}{d^2} \left(\left(1 - \frac{1}{\epsilon^2} \right) r^2 + \frac{2d}{\epsilon^2} r \right) - 1$$

so:

$$\frac{1}{d} = \frac{x^2 m^2 M G}{L^2} \quad - (21)$$

$$\boxed{d = \frac{L^2}{x^2 m^2 M G}} \quad - (22)$$

$$\frac{\epsilon^2}{d^2} \left(1 - \frac{1}{\epsilon^2} \right) = \frac{2m\omega^2 E}{L^2} \quad - (23)$$

$$\epsilon^2 - 1 = \frac{2d^2 m \omega^2 E}{L^2} \quad - (24)$$

$$= \frac{2L^4 m \omega^2 E}{x^4 m^4 M^2 G^2 L^2}$$

$$= \frac{2L^2 E}{x^2 m^3 M^2 G^2}$$

So

$$\epsilon = \left(1 + \frac{2L^2 E}{x^2 m^3 M^2 G^2} \right)^{1/2}$$

— (25)

It is clear that the precession constant x can be calculated classically.
