

The Convective or Lagrangian Derivative in Plane Polar Coordinates : Spin Comets and Matrix Form

In plane polar coordinates the convective derivative is given as:

$$\begin{aligned} \frac{D\bar{v}}{Dt} &= \frac{\partial \bar{v}}{\partial t} + (\bar{v} \cdot \nabla) \bar{v} = \frac{\partial \bar{v}}{\partial t} + \\ &\left(v_r \frac{\partial}{\partial r} + \frac{v_\theta}{r} \frac{\partial}{\partial \theta} + v_z \frac{\partial}{\partial z} \right) \left(v_r \underline{\mathbf{e}}_r + v_\theta \underline{\mathbf{e}}_\theta + v_z \underline{\mathbf{k}} \right) - (1) \\ &= v_r \frac{\partial}{\partial r} (v_r \underline{\mathbf{e}}_r) + \frac{v_\theta}{r} \frac{\partial}{\partial \theta} (v_r \underline{\mathbf{e}}_r) + v_z \frac{\partial}{\partial z} (v_r \underline{\mathbf{e}}_r) \\ &+ v_r \frac{\partial}{\partial r} (v_\theta \underline{\mathbf{e}}_\theta) + \frac{v_\theta}{r} \frac{\partial}{\partial \theta} (v_\theta \underline{\mathbf{e}}_\theta) + v_z \frac{\partial}{\partial z} (v_\theta \underline{\mathbf{e}}_\theta) \\ &+ v_r \frac{\partial}{\partial r} (v_z \underline{\mathbf{k}}) + \frac{v_\theta}{r} \frac{\partial}{\partial \theta} (v_z \underline{\mathbf{k}}) + v_z \frac{\partial}{\partial z} (v_z \underline{\mathbf{k}}) + \frac{\partial \bar{v}}{\partial t} \end{aligned}$$

In general the derivatives of quantities inside brackets must be worked out with Leibnitz Theorem.

In plane polar system:

$$\frac{\partial}{\partial \theta} = \frac{\partial \underline{\mathbf{e}}_r}{\partial \theta} = \frac{\partial \underline{\mathbf{e}}_r}{\partial z} = \frac{\partial \underline{\mathbf{e}}_\theta}{\partial z} = \frac{\partial \underline{\mathbf{e}}_r}{\partial z} = \frac{\partial \underline{\mathbf{k}}}{\partial z} = 0 \quad -(2)$$

and

$$\frac{\partial r}{\partial \theta} = 0. \quad -(3)$$

By construction:

$$\frac{\partial \underline{\mathbf{e}}_r}{\partial \theta} = \underline{\mathbf{e}}_\theta ; \frac{\partial \underline{\mathbf{e}}_\theta}{\partial \theta} = -\underline{\mathbf{e}}_r \quad -(4)$$

This means that :

$$\frac{D \underline{v}}{Dt} = \frac{\partial \underline{v}}{\partial t} + \left(v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} + v_z \frac{\partial v_r}{\partial z} \right) \underline{e}_r + \frac{v_\theta v_r}{r} \frac{\partial \underline{e}_r}{\partial \theta} \\ + \left(v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + v_z \frac{\partial v_\theta}{\partial z} \right) \underline{e}_\theta + \frac{v_\theta^2}{r} \frac{\partial \underline{e}_\theta}{\partial r} \\ + \left(v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z} \right) \underline{e}_z \quad - (5)$$

$$= \frac{\partial \underline{v}}{\partial t} + \begin{pmatrix} \frac{\partial v_r}{\partial r} & \frac{1}{r} \frac{\partial v_r}{\partial \theta} & \frac{\partial v_r}{\partial z} \\ \frac{\partial v_\theta}{\partial r} & \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} & \frac{\partial v_\theta}{\partial z} \\ \frac{\partial v_z}{\partial r} & \frac{\partial v_z}{\partial \theta} & \frac{\partial v_z}{\partial z} \end{pmatrix} + \begin{pmatrix} 0 & -\frac{v_\theta}{r} & 0 \\ \frac{v_\theta}{r} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} v_r \\ v_\theta \\ v_z \end{pmatrix}$$

ii cylindrical polar coordinates.

In the plane polar coordinates :

$$\frac{D \underline{v}}{Dt} = \frac{\partial \underline{v}}{\partial t} + \left(\begin{bmatrix} \frac{\partial v_r}{\partial r} & \frac{1}{r} \frac{\partial v_r}{\partial \theta} \\ \frac{\partial v_\theta}{\partial r} & \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} \end{bmatrix} + \begin{bmatrix} 0 & -\frac{v_\theta}{r} \\ \frac{v_\theta}{r} & 0 \end{bmatrix} \right) \begin{bmatrix} v_r \\ v_\theta \end{bmatrix} \quad - (6)$$

The spin conversion matrix of the plane polar system of coordinates therefore :

$$\omega_{ab}^a = \begin{bmatrix} \frac{\partial v_r}{\partial r} & \frac{1}{r} \frac{\partial v_r}{\partial \theta} \\ \frac{\partial v_\theta}{\partial r} & \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} \end{bmatrix} + \begin{bmatrix} 0 & -\frac{v_\theta}{r} \\ \frac{v_\theta}{r} & 0 \end{bmatrix} \quad - (7)$$

) and the coriolis derivative is the Coriolis derivative:

$$\frac{Dv^a}{Dt} = \frac{\partial v^a}{\partial t} + \omega^a_{\text{ob}} v^b - (8)$$

Eq. (8) is intended to mean:

$$\frac{D}{Dt} \begin{bmatrix} v_r \\ v_\theta \end{bmatrix} = \frac{\partial}{\partial t} \begin{bmatrix} v_r \\ v_\theta \end{bmatrix} + \left(\begin{bmatrix} \frac{\partial v_r}{\partial r} & \frac{1}{r} \frac{\partial v_r}{\partial \theta} \\ \frac{\partial v_\theta}{\partial r} & \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} \end{bmatrix} + \begin{bmatrix} 0 & -\frac{v_\theta}{r} \\ \frac{v_\theta}{r} & 0 \end{bmatrix} \right) \begin{bmatrix} v_r \\ v_\theta \end{bmatrix} - (9)$$

Denoting:

$$v^1 = v_r ; v^2 = v_\theta - (10)$$

$$\begin{aligned} \text{then: } \frac{Dv^1}{Dt} &= \frac{\partial v^1}{\partial t} + \left(\left(\frac{\partial v^1}{\partial r} \right) v^1 + \frac{1}{r} \left(\frac{\partial v^1}{\partial \theta} \right) v^2 \right) - \frac{v^2}{r} \\ &= \frac{\partial v^1}{\partial t} + \left(\left(\frac{\partial v^1}{\partial r} \right) v^1 + \left(\frac{\partial v^1}{\partial \theta} - \frac{v^2}{r} \right) v^2 \right) - (11) \\ &= \frac{\partial v^1}{\partial t} + \omega^1_{01} v^1 + \omega^1_{02} v^2 \end{aligned}$$

So

$$\omega^1_{01} = \frac{\partial v^1}{\partial r} = \frac{\partial v_r}{\partial r} - (12)$$

$$\omega^1_{02} = \frac{1}{r} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta}{r} - (13)$$

$$\omega^2_{01} = \frac{\partial v_\theta}{\partial r} + \frac{v_r}{r} - (14)$$

$$\omega^2_{02} = \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} - (15)$$

similarly:

+) Eqs. (14) and (15) are worked out from:

$$\frac{D\mathbf{v}^2}{Dt} = \frac{\partial \mathbf{v}^2}{\partial t} + \omega^2_{01} v^2 + \omega^2_{02} v^2 - (16)$$

i.e.

$$\begin{aligned} \frac{D\mathbf{v}_\theta}{Dt} &= \frac{\partial \mathbf{v}_\theta}{\partial t} + \omega^2_{01} v_r + \omega^2_{02} v_\theta \\ &= \frac{\partial \mathbf{v}_\theta}{\partial t} + \left(\frac{\partial \mathbf{v}_\theta}{\partial r} + \frac{v_\theta}{r} \right) v_r + \left(\frac{1}{r} \frac{\partial \mathbf{v}_\theta}{\partial \theta} \right) v_\theta \end{aligned} - (17)$$

so eqs. (14) and (15) follow, Q.E.D.

Now we:

$$v_r = \dot{r}, \quad v_\theta = r\dot{\theta} - (18)$$

to find out:

$$\begin{bmatrix} 0 & -v_\theta/r \\ v_\theta/r & 0 \end{bmatrix} = \begin{bmatrix} 0 & -\dot{\theta} \\ \dot{r} & 0 \end{bmatrix} - (19)$$

i) the spin connection that define the angular velocity:

$$\dot{\theta} = \frac{d\theta}{dt}. - (20)$$

This means that the plane polar coordinate system is a moving frame, technically a rotating frame therefore the spin connection components are:

$$\omega^1_{01} = \dot{r}/r - (21)$$

$$\omega^1_{02} = \frac{1}{r} \frac{dr}{d\theta} - \dot{\theta} - (22)$$

$$\omega^2_{01} = \frac{\partial(r\dot{\theta})}{\partial r} + \dot{\theta} - (23)$$

$$\omega^2_{02} = \frac{1}{r} \frac{\partial(r\dot{\theta})}{\partial \theta} - (24)$$